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## Solutions to the Exercises of Tutorial 08

1. Proof. Since u satisfies the mean value formula, we have

$$u(x) = \frac{1}{\pi r^2} \int_{B_r(x)} u(y) dy,$$

for any  $B_r(x) \subset \Omega$ .

We employ the argument of contradiction and assume that there exists a point  $x_0 \in \Omega$ , such that  $\Delta u(x_0) \neq 0$ . Without loss of generality we assume that  $\Delta u(x_0) > 0$ . From the continuity of u we conclude that there exists a ball  $B_{\varepsilon}(x_0)$  such that

$$\Delta u(x) > 0$$
, for all  $x \in B_{\varepsilon}(x_0)$ , (1)

here  $\varepsilon$  is a small positive number.

Since we assume that  $u \in C_2(\Omega)$ , we can apply the Green formula to get

$$\int_{B_{\varepsilon}(x_0)} \Delta u(x) dx = \int_{\partial B_{\varepsilon}(x_0)} \nabla u(x) \cdot n dS_x$$

$$= \int_0^{2\pi} \frac{\partial}{\partial r} u(r\omega) d\theta$$

$$= \frac{\partial}{\partial r} \int_0^{2\pi} u(r\omega) d\theta$$

$$= \frac{\partial}{\partial r} u(x_0) = 0.$$

However, the integral  $\int_{B_{\varepsilon}(x_0)} \Delta u(x) dx$  is positive by (1).

Thus we arrive at a contradiction. Whence  $\Delta u(x) = 0$  for all  $x \in \Omega$ , i.e. u is harmonic.

**2.** Proof. Firstly we compute  $\Delta f(u)$  for  $C^2$ -functions  $f: \mathbb{R} \to \mathbb{R}$  and  $u: \mathbb{R}^n \to \mathbb{R}$ . There holds

$$\Delta(f(u)) = f'(u)\Delta u + f''(u)|\nabla u|^2.$$

i) Thus choosing  $f(u) = u^2$  we then have  $\Delta(u^2) = 2(u\Delta u + |\nabla u|^2)$ . Since u is a harmonic function, one has

$$\Delta(u^2) = 2|\nabla u|^2 \ge 0.$$

Now we choose a ball  $B \subset \in \mathbb{R}^n$ . It is easy to find a solution v to

$$\Delta v = 0$$
, in B,

and  $v|_{\partial B} = (u^2)|_{\partial B}$ . We need only to prove that  $v \leq u$  in B. To prove this we let  $w = u^2$  and F = w - v. Then F satisfies

$$\Delta F = g$$
, in  $B$ ,

and  $F|_{\partial B}=0$ , where  $g=2|\nabla u|^2\geq 0$ . Applying the maximum principle we obtain

$$F \leq 0$$
.

Therefore  $u^2$  is sub-harmonic.

- ii) For any convex function  $f \in C_2(\mathbb{R}, \mathbb{R})$ , by a slightly different argument we prove that the composite function f(u) is sub-harmonic.
- **3.** Proof. The proof consists of three steps. We first construct a smooth approximation, say  $|x|_{\varepsilon}$ , of |x|, then prove this approximate function is sub-harmonic. Finally pass the approximate function to limit as  $\varepsilon \to 0$ .
  - i) We define

$$|x|_{\varepsilon} = \sqrt{|x|^2 + \varepsilon^2}$$
.

It is easy to show that  $|x| + \sqrt{|x|^2 + \varepsilon^2} \ge \varepsilon$  for all x. Thus

$$||x|_{\varepsilon} - |x|| = \frac{\varepsilon^2}{|x|_{\varepsilon} + |x|} \le \varepsilon \to 0$$

as  $\varepsilon \to 0$ .

ii) Straightforward computations yield

$$\partial_{x_i}(|x|_{\varepsilon}) = \frac{x_i}{|x|_{\varepsilon}},$$

and

$$\partial_{x_i}^2(|x|_{\varepsilon}) = \frac{|x|^2 + \varepsilon^2 - x_i^2}{|x|_{\varepsilon}^3},$$

from which we obtain

$$\Delta(|x|_{\varepsilon}) = \frac{(n-1)|x|^2 + n\varepsilon^2}{|x|_{\varepsilon}^3} =: f_{\varepsilon}(x).$$

Choose any ball  $B \subset \mathbb{R}^n$ . Suppose that  $v^{\varepsilon}$  is a solution to

$$\Delta v^{\varepsilon} = 0$$
, in B.

and  $v^{\varepsilon}|_{\partial B} = (|x|_{\varepsilon})|_{\partial B}$ . From the lecture notes we know such solution exists. Let  $w_{\varepsilon} = |x|_{\varepsilon} - v^{\varepsilon}$ , then w satisfies

$$\Delta w = f$$
, in  $B$ ,

and  $w_{\varepsilon}|_{\partial B}=0$ . Note that  $f_{\varepsilon}\geq 0$ . Applying the maximum principle yields

$$w \leq 0$$

whence  $|x|_{\varepsilon}$  is sub-harmonic.

iii)  $|x|_{\varepsilon}$  converges to |x| uniformly. Note that

$$|x| \leq |x|_{\varepsilon} \leq v^{\varepsilon}$$
, in B

and  $(v^{\varepsilon})|_{\partial B} = |x|_{\varepsilon}^{\varepsilon}|_{\partial B} \to (|x|)|_{\partial B}$  as  $\varepsilon \to 0$ . We also have  $v^{\varepsilon} \to v$  in  $C(\bar{B})$  which is harmonic. Thus |x| is sub-harmonic. This argument indicates that Theorem 5.12 in the lecture notes is still true for the case that we have infinitely many sub-solutions.

**4.** i) Define  $w \equiv 0$ . It is easy to see that w is a super-solution to the equation

$$\lambda v + \Delta v = 0$$
, in  $\Omega$ ,

with  $\lambda \leq 0$  satisfying the boundary condition  $w|_{\partial\Omega} = 0$ .

- ii) There are two cases according to  $\lambda = 0$  or  $\lambda < 0$ .
- a) We assume firstly that  $\lambda < 0$ . Define a function in the form:

$$u_1(x) = -M \left(1 - e^{-\mu d(x)}\right)$$

Here M,  $\mu > 0$  are positive constants to be determined later on, and d = d(x) defined  $d(x) = \operatorname{dist}(x, \partial\Omega)$  for any  $x \in \Omega$ , is the distance function which is Lipschitz continuous (this can be proven by using the triangle inequality easily). Furthermore, since the boundary  $\partial\Omega$  is  $C^2$ , we have that d(x) is  $C^2$  too for x near  $\partial\Omega$ ,  $u_1$  is also  $C^2$ . Now we calculate

$$\partial_i u_1 = -\mu M e^{-\mu d} \partial_i d,$$

and

$$\partial_i^2 u_1 = (\mu(\partial_i d)^2 - \partial_i^2 d) M \mu e^{-\mu d}$$

Therefore, we obtain for x near  $\partial\Omega$ , say  $x\in\Omega_{\gamma}=\left\{x\in\Omega\mid d(x)<\frac{1}{\gamma}\right\}$ , that

$$\lambda u_1 + \Delta u_1 = -\lambda M + \left(\lambda + \mu^2 \sum_{i=1}^3 (\partial_i d)^2 - \mu \Delta d\right) M e^{-\mu d}$$

$$\geq -\lambda M + \frac{1}{2} \lambda M$$

$$= -\frac{1}{2} \lambda M > 0,$$

provided that  $x \in \Omega_{\gamma}$  and  $\gamma = \frac{1}{\mu}$ , where  $\mu$  is sufficiently small from which it follows that  $e^{-\mu d}$  approaches to 1 (In fact, we have  $e^{-1} \le e^{-\mu d} \le 1$  for all  $x \in \Omega_{\gamma}$ ), also the

continuity of  $\Delta d$  and  $\partial_i d$  over  $\bar{\Omega}_{\gamma}$  which implies that they are uniformly bounded on  $\bar{\Omega}_{\gamma}$ , namely there exists a constant K such that  $|\partial_i d| + |\Delta d| \leq K$ .

The above computations motivate us to construct a sub-solution as follows: Choose the two positive constants M, C such that  $C \ll M$ . Then we have

$$\underline{u}(x) = u_1(x)$$

for x which closes sufficiently to the boundary of  $\Omega$ , otherwise

$$u(x) = -C$$
.

Using the property that  $\max\{f,g\}$  is a sub-solution if f,g are sub-solutions (see also Theorem 5.12 in the lecture notes), one can easily prove that  $\underline{u}$  is a sub-solution. to the Helmholtz equation, satisfying  $\underline{u}|_{\partial\Omega}=0$ . It remains to show that  $u_1$  is a sub-solution.

Since d may be not differentiable at x which is not near the boundary, we employ an approximate procedure. One way is to apply the Weierstrass theorem which says that a continuous function can be approximated by a polynomial sequence. Another is to use the convolution introduced in the lecture. Here we use the latter one. Taking a function  $\varphi_{\varepsilon} \in C_{\infty}$  such that  $\int_{\mathbb{R}^n} \varphi_{\varepsilon} dx = 1$ , the support of  $\varphi_{\varepsilon}(x)$  is  $B_{\varepsilon}(0)$ . It is not difficult to prove that

$$f * \varphi_{\varepsilon} \in C_{\infty}(\mathbb{R}^n),$$

if  $f \in L_{1,loc}(\mathbb{R}^n)$ . Applying this to function d which is Lipschitz continuous, hence  $d \in L_1(\Omega)$ , so  $d * \varphi_{\varepsilon} \in C_{\infty}(\Omega)$ . Moreover, we have

$$d * \varphi_{\varepsilon} \to d$$
, in  $C(\Omega)$ .

We now define a new function  $u_1^{\varepsilon}$  by

$$u_1^{\varepsilon}(x) = -M \left(1 - e^{-\mu d^{\varepsilon}(x)}\right)$$

Similar the computations carried out for  $u_1$ , and choosing  $\mu$  suitably small (however maybe depends on  $\varepsilon$  in this case) yield

$$\begin{split} \lambda u_1^\varepsilon + \Delta u_1^\varepsilon &= -\lambda M + \left(\lambda + \mu^2 \sum_{i=1}^3 (\partial_i d^\varepsilon)^2 - \mu \Delta d^\varepsilon \right) M e^{-\mu d^\varepsilon} \\ &\geq -\lambda M + \frac{1}{2} \lambda M \\ &= -\frac{1}{2} \lambda M > 0, \end{split}$$

for any fixed  $\varepsilon$ . From this one can conclude that  $u_1^{\varepsilon}$  is a sub-solution to Helmholtz equation satisfying suitable boundary condition. Using again Theorem 5.12 in a slightly general version (i.e. for infinitely many sub-solutions), we then conclude

the limit,  $u_1$ , of  $u_1^{\varepsilon}$ , is also a sub-solution to the Helmholtz equation satisfying the zero boundary condition.

b) We now consider the case that  $\lambda=0$ . For this problem, the above technique does not work. However it is valid for a new function  $v^{\varepsilon}$  which is defined by  $v^{\varepsilon}=e^{\varepsilon x_1}u$  with a small parameter  $\varepsilon$ . For simplicity we omit the upper-script  $\varepsilon$  till we investigate its limit. Then one has

$$\partial_1 v = \varepsilon e^{\varepsilon x_1} u + e^{\varepsilon x_1} \partial_1 u, \quad \partial_i v = e^{\varepsilon x_1} \partial_i u, \text{ for } i \neq 1,$$

and

$$\partial_1^2 v = \varepsilon^2 e^{\varepsilon x_1} u + 2\varepsilon e^{\varepsilon x_1} \partial_1 u + e^{\varepsilon x_1} \partial_1^2 u, \quad \partial_i^2 v = e^{\varepsilon x_1} \partial_i^2 u, \text{ for } i \neq 1,$$

Thus

$$\Delta v + 2\varepsilon \partial_1 v - 3\varepsilon^2 v = 0.$$

There is a term in the equation, which involves the first order derivative  $\partial_1 v$ . However, from the argument in a) we see that this is not an essential difficulty. Letting  $\lambda = -3\varepsilon^2$ , we see that this equation satisfies the requirement of a), i.e.  $\lambda < 0$ .

Therefore we can construct a sub-solution satisfying  $v^{\varepsilon}|_{\partial\Omega} = 0$ , in a similar manner as in a), for every fixed  $\varepsilon$ . Then letting  $\varepsilon \to 0$  we get the limit function u satisfying  $u|_{\partial\Omega} = 0$  and u is a sub-solution to the Helmholtz equation.